

SUBLINEAR INTENSITY RESPONSE OF CERIUM-DOPED YTTTRIUM ALUMINIUM GARNET SCREEN WITH CHARGE

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ABSTRACT

- APS-U swap-out injection necessitates ~ 17 nC electron bunches at 6 GeV.
- Diagnostic imaging screens are envisaged in the BTS.
- Important to determine whether the response of these screens to charge is linear.
- In the present work, we examine the effect of sublinear intensity quenching of a Cerium-doped Yttrium-Aluminium-Garnet scintillator screen.
- At 7 GeV, charge density $\rho \leq 10$ fC μm^{-2} , approximately 10 % reduction of the imaging intensity due to scintillator quenching.

MOTIVATION

- Saturation of scintillator screens has represented a challenge for beam imaging at many facilities – in particular linacs – e.g. recently at Euro-XFEL [1].
- At the APS, prior work on scintillator linearity included experiments using the electron linac [2, 3].
- With high-charge bunches through the BTS transport line for the Advanced Photon Source Upgrade (APS-U), will scintillator linearity with charge be a significant detrimental effect?

SCINTILLATOR QUENCHING

- Scintillator quenching occurs when charge density of incident beam depletes the vacancies in the crystal, and the crystal does not produce light output at a rate proportional to the input charge density [4–6].
- We consider limits to quenching of the scintillator along the theory of Birks [7].
- Results in approximate upper charge density limits of 16 fC μm^{-2} for LYSO scintillators [8], and 20 fC μm^{-2} for YAG:Ce scintillators [9].
- Quenching is possible in Chromox ($\text{Al}_2\text{O}_3:\text{Cr}$), however contemporary applications of Chromox scintillators for imaging are typically proton rather than electron beams.
- Optical Transition Radiation (OTR) has no quenching limit: the limit is probably the damage threshold of the material surface. In practice, if an electron bunch is short (\sim tens of fs duration), the practical limit for OTR is probably the presence of Coherent Optical Transition Radiation.
- Even for bunches of \sim ps duration, COTR will occur when there is microbunching, or if there is a narrow current spike.

METHOD

- In the present work, we use a 7 GeV electron beam coming from the booster, imaged using the fluorescent screen BTS:FS3 [10].
- Quenching of the scintillator reduces the light output of the scintillator at locations on the screen with highest charge density.
 - In effect, this results in fitting the ‘tails’ of the distribution, and essentially it appears that the image of the beam on a scintillator is larger than the rms electron beam size.
- We evaluate the electron beam size as a bivariate Gaussian distribution, in order to quantify the areal charge density.
 - The equation of a bivariate Gaussian distribution in coordinates x_i , with means μ_i , standard deviations σ_i and is given by [2008]:
- Hence an electron beam with a profile that is Gaussian in two dimensions with root mean square beam sizes σ_1, σ_2 , we can describe the peak electron charge density ρ by:

$$p = \frac{1}{2\pi\sigma_1\sigma_2} \left(-\frac{1}{2\sigma_1^2}(x_1 - \mu_1)^2 - \frac{1}{2\sigma_2^2}(x_2 - \mu_2)^2 \right)$$

$$\rho = \frac{q}{2\pi\sigma_1\sigma_2}$$

RESULTS – REGULAR BEAM

Figure 1: Images of the beam with charge. (a) 0.52 nC. (b) 1.06 nC. (c) 3.2 nC. (d) 4.6 nC.

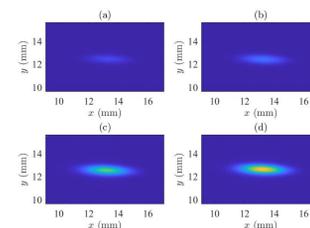


Figure 2: Fitted electron beam sizes as a function of electron beam charge.

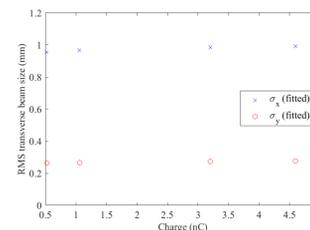


Figure 3: Fitted peak intensity as a function of electron beam charge.

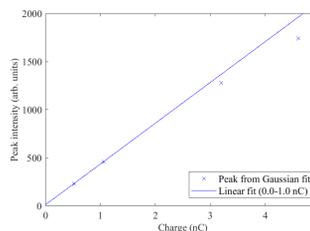
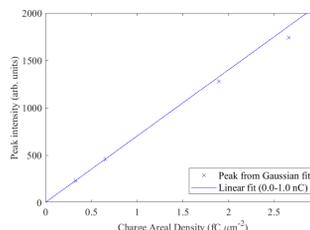


Figure 4: Intensity as a function of peak charge areal density.



RESULTS – FOCUSED BEAM

Figure 5: Images of the beam as a function of charge. (a) 0.070 nC. (b) 0.15 nC. (c) 0.50 nC. (d) 1.0 nC. (e) 1.5 nC. (f) 2.0 nC.

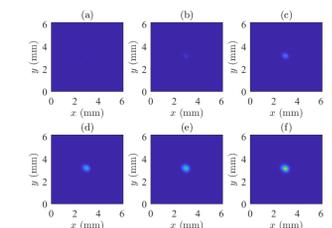


Figure 6: Fitted electron beam sizes as a function of electron beam charge.

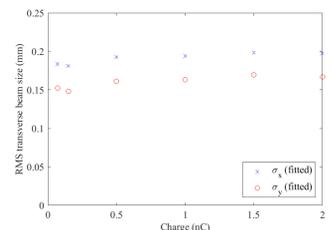


Figure 7: Fitted peak intensity as a function of electron beam charge.

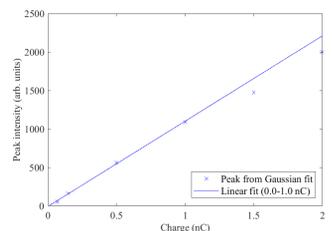
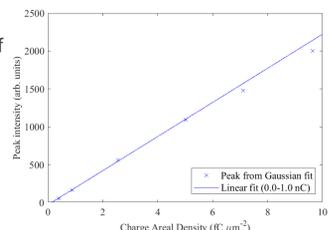


Figure 8: Intensity as a function of peak charge areal density.



DISCUSSION

- At the charge densities investigated ($\rho \leq 10$ fC μm^{-2}), even though we start to observe nonlinear behaviour of the scintillator, Figs. 2, 6 show negligible change to the measured beam sizes.
- However, at charge areal densities $\rho \geq 10$ fC μm^{-2} , one should anticipate that the size of the electron beam determined from the scintillation profile would become increasingly unreliable.

SUMMARY

- In conclusion, measurements of the electron beam on the BTS:FS3 screen were made.
- At 7 GeV beam energy leaving the booster, over the charge densities investigated ($\rho \leq 10$ fC μm^{-2}), an approximately 10 % reduction of the imaging intensity due to quenching of the scintillator was observed.

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